

**EXPERIMENTAL STUDIES ON THE OPTICAL PROPERTIES
OF
CHOLESTERIC LIQUID CRYSTALS**

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for the degree of
DOCTOR OF PHILOSOPHY

by
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DECLARATION

I hereby declare that the thesis was
composed by me independently and that
it has not formed the basis for the
award of any Degree, Diploma, Associateship,
Fellowship or other similar title.

K. A. Suresh

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Certified
S. Chandrasekar
27/1/77

CERTIFICATE

I **certify** that this **thesis** has been **composed** by Mr. K.A. **Suresh** based on **investigations** carried out by him at the **Liquid Crystals** Laboratory, **Raman** Research Institute, under my **supervision**. The **subject** matter of **this thesis** has not **previously** formed the **basis** of the award of any **degree, diploma, associateship, fellowship** or other similar **title**.



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PREFACE

Liquid crystals are states of matter that are **intermediate** between **the** crystalline solid and the isotropic liquid. **They** are strongly anisotropic in some of their **properties** (for example, they are **birefringent** and **diamagnetically** anisotropic) and at the same time they have the ability to flow. **Friedel (1922)**, after a careful study of the optical properties of the '**mesophases**', classified them into **three** types: nematic, emectic and cholesteric.

1) **The** nematic liquid arystal has long range orientational order but no translational order (**Figure 1a**). A monodomain **sample** of a nematlo is optically uniaxial, positive and strongly birefringent. The birefringence decreases gradually with rise of **temperature** and drops abruptly **to** zero at the **nematic-isotropic** transition temperature.

2) Smectia liquid **crystals** have stratified structures. Depending on the molecular arrangement within **each** layer, **smectics** are classified into **smectic** A, smectic B, smectio C, etc. **In smectic** A the long **axis** of the molecules are, on the average,

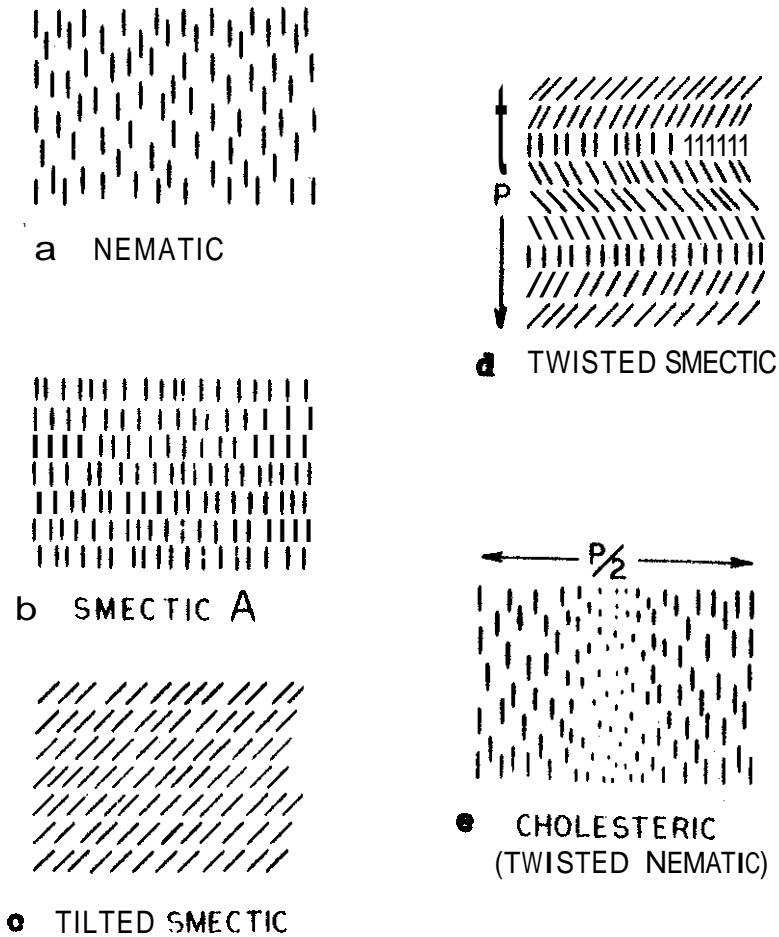


Figure 1: The molecular arrangement in various mesophases.

normal to the layers with their centres distributed at random within each layer (**Figure 1b**). **Smectic B** resembles **smectic A** except that the molecules are hexagonally close packed in each layer. **Smectic C** is a tilted form of **smectic A** (**Figure 1c**). **Recently** it has been shown **that** **smectic C** may also occur in a twisted form, **i.e.**, with a twist imposed on the structure about **an** axis normal to the layers (**Figure 1d**). **The** other **smectic modifications, e.g., D, E, F, etc.** have more **complicated** molecular structures which are not yet fully **elucidated**.

3) Cholesteric liquid crystals are also nematic type of liquid crystals, except that they are composed of optically active molecules. Thus, the structure has a spontaneous **helical** twist superimposed about **an** axis perpendicular to the preferred molecular direction (**Figure 1e**). For the purpose of optical calculations the structure may also be described as a pile of **very** thin **birefringent** layers, the principal **axes** of the successive layers turned through a small angle. This spiral structure gives rise to some remarkable optical properties. **When** white light is incident on the surface of a plane texture preparation (in which the helical **axis** is normal to the **glass** surfaces),

selective reflection takes place over a small region of the spectrum, the wavelength of maximum reflection varying with angle of incidence in accordance with **Bragg's** law. At normal incidence, the reflected light is circularly **polarised**: one circular component **is** totally reflected over a **spectral range** of some 0.02 μm while the other passes through unchanged. **Further** in contrast to the case of reflection from **normal** dielectrics, the reflected wave has the same sense of **circular** polarisation **as** that of the incident wave. Along the helical axis, the **medium** exhibits optical rotatory power of the order of **several** thousands of **degrees** per **millimetre**. In the **neighbourhood** of the region of reflection, the rotatory dispersion is anomalous and the sign of the **rotation** opposite on opposite sides of the reflection band.

These properties are generally well understood theoretically and have been extensively studied **experimentally**. However, the **experimental** and theoretical studies have not been fully explored in certain special cases. **This** thesis **describes** investigations **carried** out on **such** problems. The work has led to some **significant** new results.

The first two **chapters** deal with studies **in** the

regime of large **pitch**, much larger than the wavelength of light. These studies **have** established conclusively that the optical properties in this regime cannot be described adequately by a direct application of the **standard de Vries** equation as has been done by **previous** workers (**Baessler et al.** 1969, **Teucher et al.** 1971). The **experiments** were carried out on a mixture of right **handed** and **left** handed cholesterics. Such a mixture, referred to as a compensated mixture, adopts the helical **structure** of a cholesteric but the **pitch** is sensitive to **composition** and temperature. For a given composition, there is a temperature (say T_N) at which the **pitch** becomes infinity, **i.e.**, structure goes over to the nematic state (**Friedel**, 1922). At temperatures above and below T_N the structure is cholesteric but of opposite handedness. The particular compensated **cholesteric** used in the present work was a **1.6:1** by weight mixture of cholesteryl chloride and **cholesteryl** myristate (CM), which has a nematic temperature $T_N \approx 48^\circ\text{C}$. Some measurements of the rotatory power as a function of **temperature/pitch** for this mixture (but of composition 1.75: 1) have been reported previously by Baessler **et al.** (1969), who interpreted the observations on the **basis** of the de Vries equation,

$$\rho = -\frac{2\pi}{P} \frac{a^2}{8\lambda'^2(1-\lambda'^2)} \quad (1)$$

where ρ is the optical rotatory power, P is the pitch of the helix, $a = (\epsilon_1 - \epsilon_2)/2\epsilon$, ϵ_1 and ϵ_2 are the principal dielectric constants of the **untwisted** structure, $\epsilon = \frac{1}{2}(\epsilon_1 + \epsilon_2)$, $\lambda' = \frac{\lambda}{P\epsilon^{\frac{1}{2}}}$ and λ is the wavelength in **vacuum**.

In chapter I we present systematic measurements of the rotatory power as a function of temperature/pitch for samples of **thickness** 3, 6.4 and 12.7 μm . These measurements were carried out using a slightly **modified Winkel-Zeiss (Göttingen) Polarimeter**. The experimental curves are **shown** in figure 2. It can be seen that the rotatory power (**i.e.**, the optical rotation per unit thickness) as well as **its** variation with **temperature/pitch** is a function of the thickness of the sample. It is evident **that** this **behaviour cannot** be explained in terms of equation (1). According to this equation the rotatory power ρ should **approach infinity** as temperature T tends to T_N (or pitch tends to infinity) and should be independent of sample thickness, whereas experimentally ρ tends to zero in the vicinity of T_N and is very much dependent on the sample thickness. The experimental results are

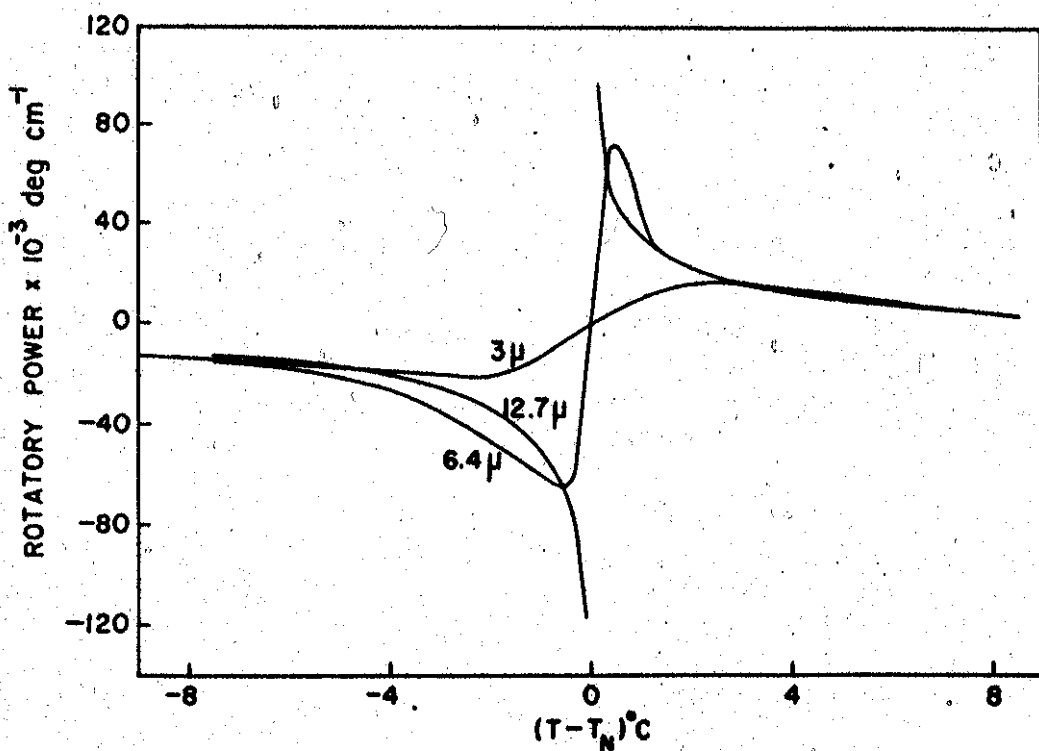


Figure 2: Experimental rotatory power as a function of temperature for a 1.6:1 by weight mixture of cholesteryl chloride and cholesteryl myristate. Sample thicknesses: $\approx 3 \mu\text{m}$, $6.4 \mu\text{m}$ and $12.7 \mu\text{m}$.

discussed in the light of the rigorous theory of light propagation in oholesterics of **large pitch** developed recently by **Kini and others** (see **Kini, 1977**). It is shown that the observed behaviour is **in conformity** with the predictions of this exact theory.

Chapter **II** deals with an absorbing **compensated** mixture. Saeva and **Wysocki** (1971) demonstrated that when linearly dihydroio molecules are dissolved in a **cholesteric liquid crystal** the medium exhibits **circular dichroism** in the region of the absorption band of the solute molecules. In the present work, we have **carried** out detailed measurements of the circular dichroism as a function of **temperature/pitch** of the CM **mixture** in which was **dissolved 1.38%** by weight of linearly dichroic **β -carotene**. The circular dichroism was determined by accurate measurements of the intensity of right and left-circular polarised light transmitted by the sample, using a photomultiplier **and** a lock-in-amplifier. As is to be expected the **circular dichroism** changes sign on crossing **T_N** in a manner **somewhat** similar to the behaviour of the rotatory power of a non-absorbing sample. However, **measurements on samples** of thicknesses 4, 6.3, 8, 9.5 and 12.7 microns yielded a most unexpected **result**: the dihydroio power (**i.e.**,

the **circular dichroism per unit thickness** of the sample) was found to decrease with increasing sample thickness (see Figure 3). These **features** cannot be interpreted on the basis of Sackmann and Voss (1972) equations obtained in the **limits** $P \ll \frac{\lambda}{\Delta\mu}$ and $P \gg \frac{\lambda}{\Delta\mu}$ (λ is the wavelength of light in **vacuum** and $\Delta\mu$ is the birefringence). It is **shown** that the observed features can be properly accounted for on the **basis** of an **extension** of the rigorous theory to absorbing systems (Kini, 1977).

Chapters III and IV are concerned with the **optical** properties in the regime of wavelengths comparable to the **pitch**. The exact theory of light propagation **along** the helical **axis** of a cholesteric was developed by Mauguin (1911), Oseen (1933) and de Vries (1951) and has been **discussed** in **various forms** by other authors (Kats 1971, Nityananda 1973, Joly 1972, Aihara and Inaba 1971, Marathay 1971). An alternate approach to the problem was put forward by Chandrasekhar and Srinivasa Rao (1968). This theory is **closely** analogous to Darwin's **dynamical** theory of X-ray diffraction **from** crystals (1914). In chapter III we derive this theory in a complete form avoiding certain ^{minor} **inconsistencies** that were present

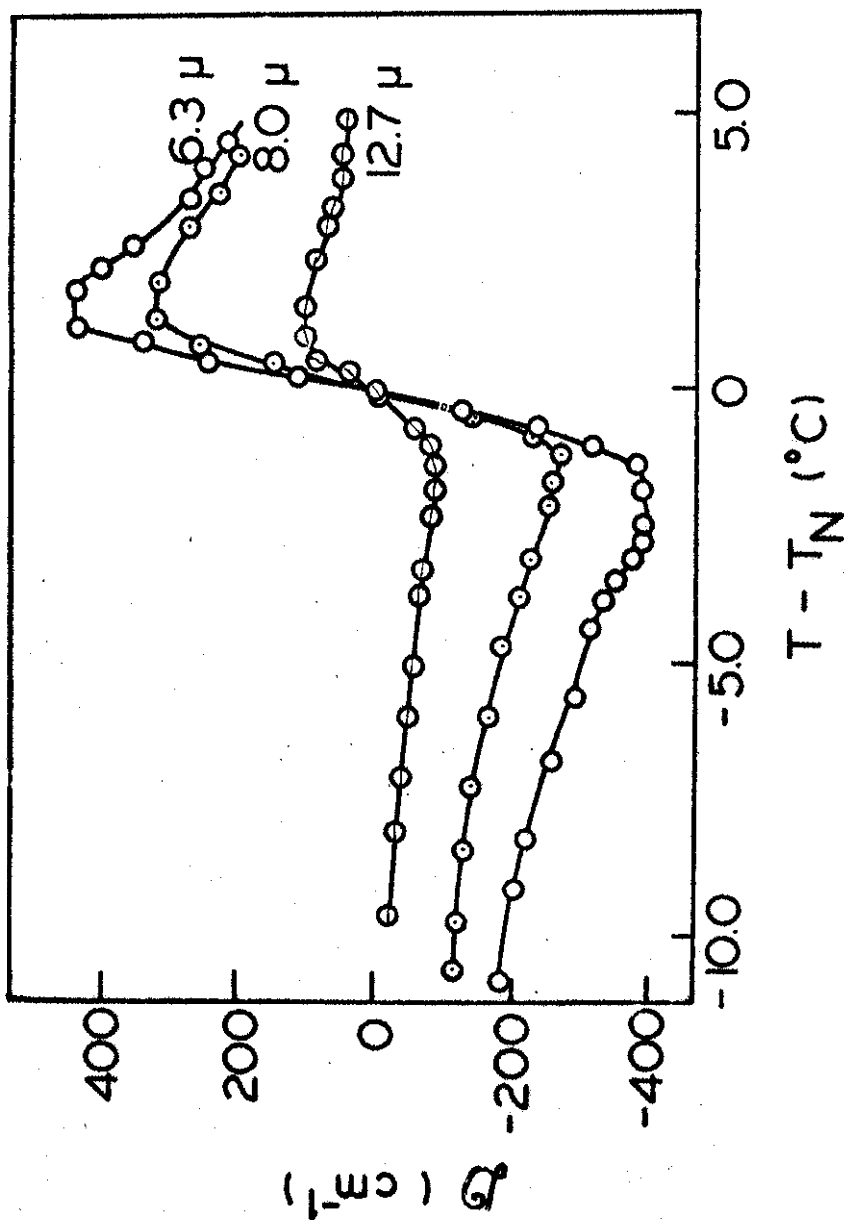


Figure 3: Experimental values of dichroic power versus temperature for different thicknesses of the sample in a 1.64:1 by weight mixture of cholesteryl chloride and cholesteryl myristate containing 1.48% by weight of β -carotene.

in the original formulation. We also give detailed **numerical** calculations of the reflection coefficient, rotatory power, etc. for the semi-infinite medium as well as for finite film and compare them with the values computed from the **exact electromagnetic** theory. It is shown that in the **neighbourhood** of the reflecting region the simple dynamical approach is quite adequate for all practical **calculations**.

In Chapter IV we discuss an entirely **new** phenomenon, namely, the optical analogue of the Borrmann **effect** in **cholesterics**. The Borrmann **effect** is the anomalous increase in the X-ray **intensity** (over **and** above that due to normal absorption) when an absorbing crystal is set for Bragg reflection (Borrmann 1941). **This** is a consequence, of the dynamical theory of X-ray diffraction (Ewald 1965). We **report** here the discovery of a similar effect in absorbing **cholesterics** **in** the vicinity of the **reflection** band? The chapter begins with a **discussion** of the **dynamical theory** (which was treated in the previous chapter) taking into account the effect of absorption. It is shown that under suitable circumstances there is an enhancement of the transmitted intensity for left (or right) circularly **polarised** light, if the structure is left (or. right)

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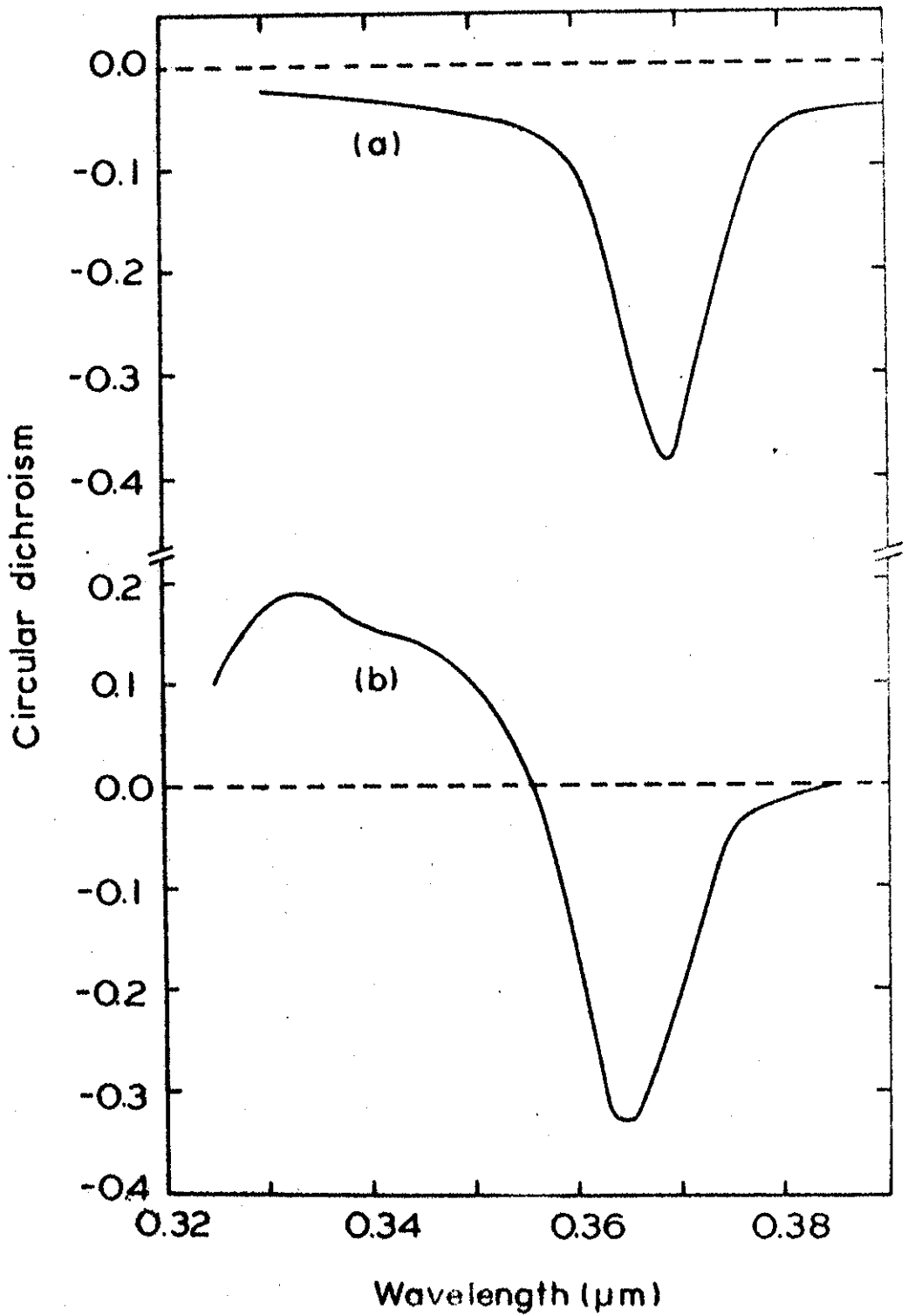
The exact electromagnetic theory of the effect was **first** worked out by Nityananda and Kini, 1973.

handed. **This is** indeed the analogue of the **Borrmann effect**. Physioally the origin of **this** effect **can** be explained on the basis of **dynamical** theory **as** follows. Sn a left handed **cholesteric** standing wares are **set up inside** the medium due to the **interference** between the **primary** and the reflected left **circularly polarised** waves, **The** phase of the **Bragg** refleoted wave with respect to the primary wave varies from π to 0 **as** λ **Increases** **from** the **shorter** wavelength edge of **the** refleotion band **to** the longer **wavelength** edge. Henoe the **electric** vector in the **medium** makes an angle $\pi/2$ **with** respect **to** the **director** on the shorter wavelength side, and **is** along the director on the longer wavelength side. Sf the medium were to **consist** of linearly **dichroic** moleoules aligned parallel to the **local** direotor **then** the left oiroular wave **suffers** **minimum** absorption at the shorter wavelength edge and maximum absorption at the longer wavelength edge of the **reflection** band. **This** results in **anomalous** increase in the transmitted **intensity** of the left oiroular wave on the ahorter wavelength side which **is** over and above the normal **attenuation** due to mean abeorption. The **existenoe**

of this effect has been experimentally demonstrated in **cholesteryl nonanoate** containing small **quantities** of **p-azoxyanisole** or **n-p-methoxybenzylidene-p-phenyl azoaniline** (which are **linearly dichroic** molecules). When the reflection band is adjusted to overlap with the **strongly dichroic** band of the solute molecule the **transmission** and circular **dichroism** (Figure 4) exhibit the **features predicted** by the theory. Also are presented a detailed experimental study of this effect in thin films of **cholesteryl nonanoate** containing **varying amounts** of **p-azoxyanisole**. It is **shown** that there is **qualitative agreement** between the **experimental results** and those **calculated** on the basis of the **dynamical** theory.

The last **chapter describes** optical studies on the twisted **smectic C (TSC) phase** and the twisted **nematic (TN or cholesteric) phase** exhibited by **p-n-heptyloxyazoxybenzene (HOAB)** mixed with **small quantities** of **cholesteryl benzoate**. The **TSC phase**, like the **TN or cholesteric phase**, also has a very high rotatory power for light propagation along the **twist axis**. **Detailed** measurements have been **carried out** of the rotatory power ρ , **pitch** P and layer birefringence ($A\mu$) **as functions** of temperature. All these **quantities** (ρ , P and $A\mu$) change **discontinuously** at the **TSC-TN phase** transition, the values being higher

Figure 4: Experimental curves of circular dichroism as a function of wavelength~(a) pure cholesteryl nonanoate (non-absorbing), (b) mixture of cholesteryl nonanoate and PAA (0.98%)(absorbing).



in the lower **temperature phase**. These studies have also led to the first **evidence** for a new **skew-oybotactio** (or tilted) form of a **cholesteric**. The **observed value** of ζ of the TN phase is found to be incompatible with **that calculated** from the **exact** theory using the measured values of $\Delta\mu$ and P when the **structure is assumed** to be of the usual cholesteric type in which the **director** is normal to the helix; the theoretical value of ζ is **much** higher, **indicating** that the **molecules** in this phase are not normal but **inclined** to the helical **axis**. This has been confirmed **directly** by **X-ray** studies on this **system**. The **diffraction** pattern reveals an inner ring (**Figure 5**) characteristic of a **skew-oybotactio** type of **cholesteric** structure in which the preferred molecular **direction** is tilted at about 39° with **respect** to the helical **axis**.

Most of the **results described** above have been reported in the **following publications** of which the present writer is a **co-author**:

1. Theory of the optical properties of **non-absorbing compensated cholesteric liquid crystals** (with S. Chandrasekhar, G.S.Ranganath, U.D. Kini), **Molecular Crystals and Liquid Crystals**, 24, 201 (1973).

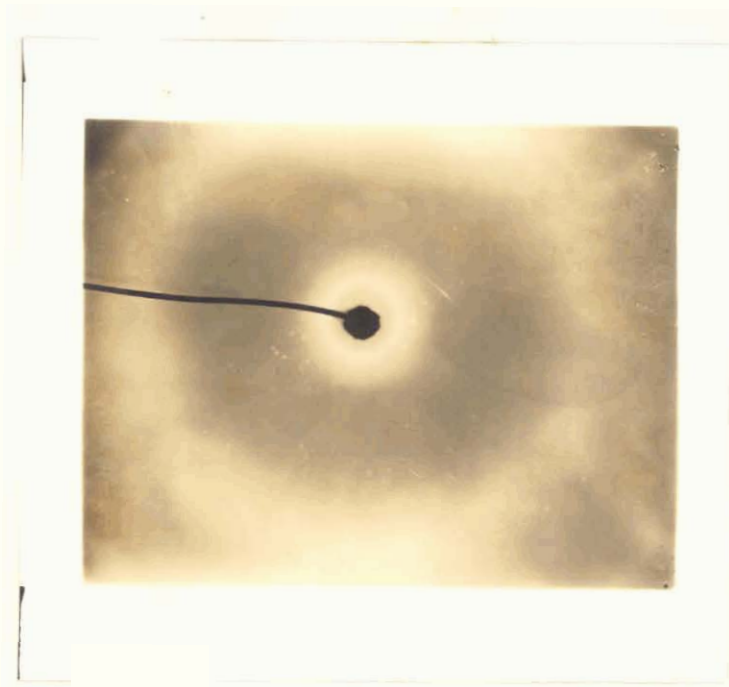


Figure 5: X-ray diffraction pattern for unaligned sample. The sample is 0.9 HOAB + 0.1 cholesteryl benzoate mixture in the twisted nematic phase.

2. **Optical** properties of mixtures of right- and left-handed cholesteric liquid crystals (with G. S. Ranganath, S. Chandrasekhar, U.D. **Kini** and S. **Ramaseshan**), *Chemical Physics Letters*, **19**, 556 (1973).
3. Circular **dichroism** in absorbing mixtures of right- and left-handed **cholesterics** (with G. S. Ranganath, **S.R.Rajagopalan** and U.D. **Kini**), **Proc. International Liquid Crystals Conference, Bangalore, December 3-8, 1973** - **Pramana Supplement 1**, p. 353
4. **Dynamical** theory of **reflection** from cholesteric liquid crystals (with S. Chandrasekhar, G. S. **Ranganath**), ^{Proc.} **International Liquid Crystals Conference, Bangalore, December 3-8, 1973** - **Pramana Supplement 1**, p. 341.
5. Anomalous transmission (Bormann effect) in absorbing cholesteric liquid crystals (with R. **Nityananda**, U.D.**Kini** and S. Chandrasekhar) - **Proc. International Liquid Crystals Conference, Bangalore, December 3-8, 1973** - **Pramana Supplement 1**, p. 325.
6. An experimental study of the anomalous transmission (Bormann effect) in absorbing cholesteric liquid crystals, *Molecular Crystals and Liquid Crystals*, **35**, 267 (1976).
7. Optical and X-ray studies on the **twisted smectic C** and **twisted nematic phases**: Evidence for a **Skew-Cybotactic** type of cholesteric **structure** (with **S.Chandrasekhar**) - *Molecular Crystals and Liquid Crystals* (in press).

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For a detailed treatment of the, structure and properties of the various mesophases, see for instance, 'Liquid Crystals' by S. Chandrasekhar, Reports on Progress in Physics, 39, 613 (1976); 'Liquid Crystals' by S. Chandrasekhar, Cambridge Honograph on Physics, Edited by M.M.Woolfson and J.M.Ziman, Cambridge University Press.